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MOTIONAL CORRELATION TIME OF DILUTE 111CD IMPURITIES IN SE-RICH-ETC(U)

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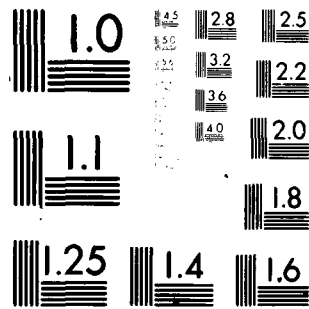
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SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER Technical Report No. 2	2. GOVT ACCESSION NO. AD-A093694	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Motional Correlation Time of Dilute $^{111}\text{Cd}$ Impurities in Se-Rich Liquid Se-Te Alloys.		5. TYPE OF REPORT & PERIOD COVERED Interim
7. AUTHOR(S) D. K. Gaskill, J. A. Gardner, K. S. Krane, K. /Krusch, and R. L. Raser (U. of Maryland Baltimore County)		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS Oregon State University Corvallis, OR 97331		8. CONTRACT OR GRANT NUMBER(s) N00014-79-C-0668
11. CONTROLLING OFFICE NAME AND ADDRESS Office of Naval Research/Chemistry Program Arlington, VA 22217		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS NR-056-721
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		12. REPORT DATE December 10, 1980
		13. NUMBER OF PAGES 5
		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report) E		
18. SUPPLEMENTARY NOTES To be published in Proceedings of the MRS Symposium on Nuclear and Electron Spectroscopies applied to Materials Science, Boston, 1980.		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) correlation time $^{111}\text{Cd}$ impurities in liquid Se-Te alloys time differential perturbed angular correlation		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) The motional correlation time, $\tau_c$ , in liquid Se and Se-rich Se-Te alloys has been investigated between 500 and 900°C using time differential perturbed angular correlations of $\gamma$ -rays from dilute $^{111}\text{Cd}$ impurities. In all alloys we find $\tau_c \propto \exp(E_0/kT)$ at low T where $E_0 = 0.36$ eV. $\tau_c$ deviates from this relation at high T. At low T, $\tau_c$ is tentatively identified as the lifetime of a Cd to host molecule bond, and at high T as the average lifetime of bonds in the host molecule.		

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SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

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OFFICE OF NAVAL RESEARCH

Contract N00014-79-C-0668

Task No. NR-056-721

TECHNICAL REPORT NO. 2

Motional Correlation Time of Dilute  $^{111}\text{Cd}$   
Impurities in Se-Rich Liquid Se-Te Alloys

by

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Prepared for Publication

in the

Proceedings of the MRS Symposium on  
Nuclear and Electron Spectroscopies Applied to  
Materials Science, Boston, 1980

Oregon State University  
Department of Physics ✓  
Corvallis, OR 97331

December 10, 1980

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# MOTIONAL CORRELATION TIME OF DILUTE $^{111}\text{Cd}$ IMPURITIES IN Se-RICH LIQUID Se-Te ALLOYS

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## ABSTRACT

The motional correlation time,  $\tau_c$ , in liquid Se and Se-rich Se-Te alloys has been investigated between 500 and 900°C using time differential perturbed angular correlations of  $\gamma$ -rays from dilute  $^{111}\text{Cd}$  impurities. In all alloys we find  $\tau_c = \exp(E_0/kT)$  at low T where  $E_0 = 0.36$  eV.  $\tau_c$  deviates from this relation at high T. At low T,  $\tau_c$  is tentatively identified as the lifetime of a Cd to host molecule bond, and at high T as the average lifetime of bonds in the host molecule.

## INTRODUCTION

The Perturbed Angular Correlation (PAC) experimental technique is not really a resonance method, but because it is a measurement of interactions between nuclei and their environment, it is informative to include it in this symposium. Physical information from PAC experiments is quite similar to that obtained from Mössbauer, NMR, and ESR hyperfine measurements.

PAC utilizes a radioactive isotope which decays via a 2-step gamma-ray cascade. The experimental work reported in this paper utilizes the isotope  $^{111}\text{In}$ , which decays to  $^{111}\text{Cd}^*$  by electron capture. The  $^{111}\text{Cd}^*$  subsequently decays by a 171 keV gamma ray ( $\gamma_1$ ) to an intermediate  $^{111}\text{Cd}$  state with lifetime  $\tau_N = 121$  nsec, which then decays to the  $^{111}\text{Cd}$  ground state by emission of a 245 keV gamma ( $\gamma_2$ ). In the absence of magnetic fields or electric field gradients at the Cd nucleus,  $\gamma_2$  would be emitted with an anisotropic probability proportional to

$$W(\theta) = [1 + A_2 P_2(\cos\theta) + A_4 P_4(\cos\theta)] \quad (1)$$

where  $\theta$  is the angle of emission of  $\gamma_2$  with respect to the direction of  $\gamma_1$ , and the nuclear angular correlation constants are given by  $A_2 = -0.180$ ,  $A_4 = 0.002$ , for this cascade. In general, however, a Cd nucleus in condensed matter will be subject to magnetic and/or electric fields due to nearby electrons and ions. Thus the nucleus will be subjected to torque during the time between emission of  $\gamma_1$  and  $\gamma_2$ , so in general the angular distribution of Eq. 1 must be modified to include the effects of reorientation of the nuclear ensemble by interaction with these fields. For a polytropic sample (i.e., one with no preferred overall orientation such as a liquid or powder) then

$$W(\theta, t) = [1 + A_2 G_2(t) P_2(\cos\theta) + A_4 G_4(t) P_4(\cos\theta)] \quad (2)$$

The functions  $G_2(t)$  and  $G_4(t)$  are normalized to unity at  $t = 0$  and are governed by the hyperfine interactions. The purpose of a PAC experiment is to measure  $G_2$  and  $G_4$  and relate them to the electronic properties of the material being investigated. Because in our case  $A_4$  is quite small, the last term in Eq. 2 contributes negligibly to  $W(\theta, t)$ , and we shall drop it. It is easy to show [1] that  $G_2(t)$  can be determined experimentally by measuring  $W(\theta, t)$  at  $\theta = 90^\circ$  and  $180^\circ$ . Then

$$A_2 G_2(t) = 2 \frac{W(180^\circ, t) - W(90^\circ, t)}{W(180^\circ, t) + 2W(90^\circ, t)} \quad (3)$$

Normally, it is possible to measure  $W(\theta, t)$  between  $t = 0$  and several times  $\tau_N$ . Details of the experimental procedure for measuring  $W(\theta, t)$  are given elsewhere. [2,3]

#### USE OF PAC TO STUDY ELECTRONIC PROPERTIES OF LIQUIDS

Both magnetic fields interacting with the nuclear magnetic moment and electric field gradients interacting with the nuclear quadrupolar moment can perturb the nucleus significantly in its intermediate state, but the latter is dominant in this experimental work, so we shall ignore magnetic fields in this discussion. We shall also discuss only the situation for which electric field gradients in the liquid fluctuate rapidly with respect to  $\tau_N$ . For such a situation, it has been shown [4] that  $G_2$  should decrease exponentially with time as

$$G_2(t) = e^{-\lambda t} \quad (4)$$

where

$$\lambda = 2.487 \tau_c \langle \nu_Q^2 \rangle \quad (5)$$

Here  $\langle \nu_Q^2 \rangle$  is the average square quadrupolar interaction frequency and  $\tau_c$  is the correlation time of electric field gradient fluctuations. The fluctuations can be caused either by physical tumbling of the molecule to which the nucleus is bound or by electronic processes which cause the chemical bonding to change rapidly. In either case  $\tau_c$  is a quantity of physical interest, but in the present experiments we have investigated a class of liquids in which  $\tau_c$  is expected to be dominated by chemical fluctuations.  $\langle \nu_Q^2 \rangle$  can be roughly estimated for these liquids and should not depend strongly on temperature. Thus our PAC measurements provide a semi-quantitative value for  $\tau_c$  and a relatively accurate measure of its temperature dependence. This information on chemical kinetics is of considerable interest and is not often measureable by other means.

#### EXPERIMENTAL RESULTS AND DISCUSSION

The work reported here on Se-rich liquid Se-Te alloys is an extension (to materials expected to have more rapid chemical bonding fluctuations) of PAC measurements on liquid Se. [2,3] For pure Se,  $\tau_c$  varies from approximately 20 psec at  $900^\circ\text{C}$  to 200 psec at  $500^\circ\text{C}$ . Over most of the range  $\tau_c$  varies with temperature as

$$\tau_c \propto e^{E_0/kT} \quad (6)$$

where  $E_0 = 0.36$  eV, but the data deviate significantly from this relation above  $800^\circ\text{C}$ . We find qualitatively similar behavior in the Se-Te alloys. The log of experimentally determined values of  $\lambda$  are shown vs  $T^{-1}$  in Fig. 1. In this figure, the right-hand ordinate shows the value of  $\tau_c$  obtained from  $\lambda$  using Eq. 5 and  $\langle v_Q^2 \rangle^{1/2} = 125$  MHz. [2] In Fig. 2 we show  $\tau_c$  for pure Se on a compressed scale and include low temperature correlation times obtained recently by Warren and Dupree [5] from  $^{77}\text{Se}$  NMR line-broadening chemical-shift anisotropy in pure liquid Se.

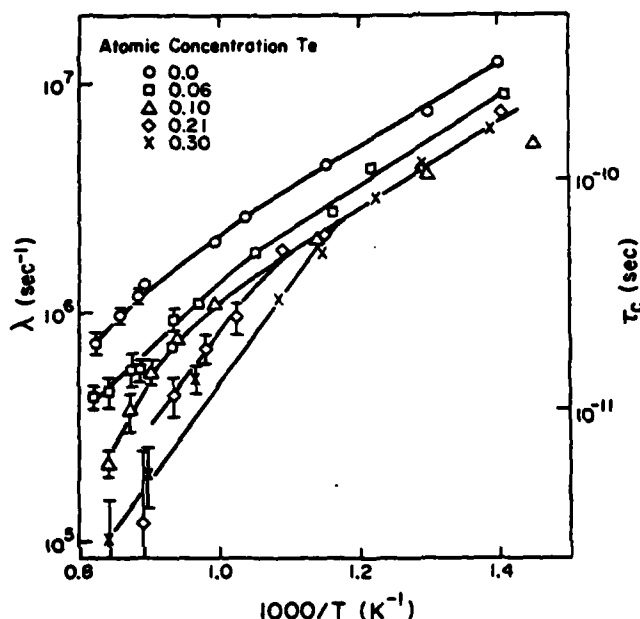


Fig. 1: Log  $\lambda$  vs  $T^{-1}$  for Cd in liquid Se-Te alloys. The right hand ordinate shows  $\tau_c$  if  $\langle v_Q^2 \rangle^{1/2} = 125$  MHz. The straight lines through the points at low  $T$  have slope  $E_0 = 0.36$  eV. The curves through the points at high temp are only a guide to the eye.

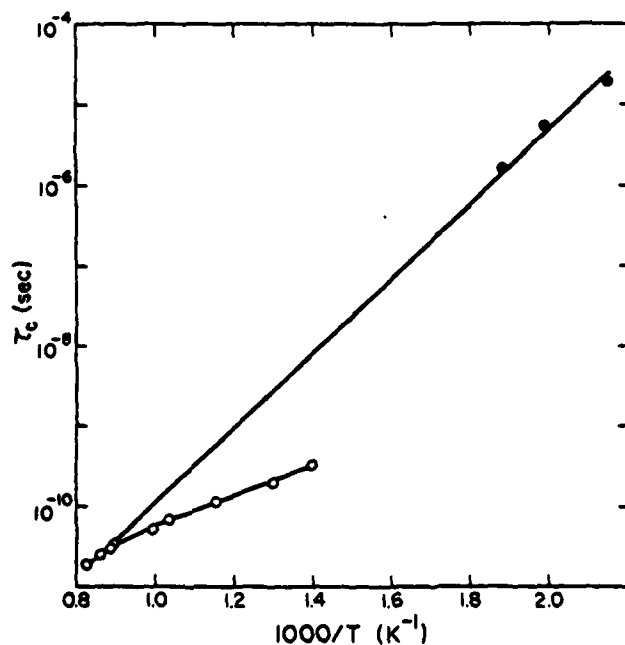


Fig. 2: Log  $\tau_c$  vs  $T^{-1}$  in liquid Se. The solid points are intrinsic correlation times measured by NMR [Ref. 5]. Open points are for Cd impurities measured by PAC, assuming  $\langle v_Q^2 \rangle^{1/2} = 125$  MHz.

Liquid Se is a weakly-bonded polymeric liquid in which the polymer size decreases with temperature due to thermal dissociation. At the highest temperatures investigated in this work, the average Se polymer size is estimated to be of order  $10^2$  atoms. [5,6] The Se-rich Se-Te alloys are qualitatively similar with somewhat weaker average bonding [7] so the polymer size and correlation times are expected to be somewhat smaller than for pure Se. Since electric field gradients at  $^{129}\text{I}$  and  $^{125}\text{Te}$  tracers are somewhat smaller in Te than in Se [8],  $\langle v_Q^2 \rangle$  may also be slightly reduced. The correlation times measured in PAC experiments will be the minimum of the polymer rotation time, the Se-Se (or in the alloys an average Se-Se, Se-Te, and Te-Te) bond lifetime, or the lifetime of the bond between the Cd tracer and the polymer molecule. In these liquids there is probably no physical difference between the first two correlation times, since the molecules almost certainly "rotate" by breaking and reforming bonds between adjacent molecules. There is of course a very important difference between this "intrinsic" correlation time and the "tracer" correlation time due to Cd bond fluctuations, and it is important to identify which of the two is the measured quantity.

The existence of two different temperature regions for  $\lambda$  which is evident in Fig. 1 suggests that in one region,  $\tau_c$  is intrinsic and is tracer-related in the other. This suggestion is further supported by the NMR measurements of  $\tau_c$  shown in Fig. 2. The NMR correlation times are intrinsic and are consistent with an extrapolation to low temperature of the pure Se high-T region but not the low-T PAC data. The implication is that the high-T PAC correlation times are intrinsic and that the low-T correlation times are extrinsic in the sense that  $\tau_c$  is dominated by Cd-polymer bond fluctuations. In principle this identification can be proven definitively by PAC measurements with other isotopes, and we are presently exploring this possibility.

It is clear from Fig. 1 that addition of small amounts of Te to Se causes an overall reduction of  $\lambda$ . Since the activation energy in the low-T "extrinsic" region is unchanged by adding Te, it is possible that at low T the change may be due to reduction of  $\langle v_Q^2 \rangle$ , and that  $\tau_c$  is not affected by alloying. This possibility can be checked by PAC measurements at sufficiently low temperature that  $\tau_c \gg \tau_N$ ,  $\langle v_Q^2 \rangle^{-1/2}$  where  $G_2(t)$  becomes independent of  $\tau_c$ , and  $\langle v_Q^2 \rangle$  can be determined directly. Although we are unlikely to achieve this condition in the liquid state, the corresponding quenched glass should be an adequate substitute. Unfortunately, our preliminary measurements at lower temperatures have been inconclusive because of experimental difficulties in sample preparation.

The PAC data in the interesting high-T region are relatively inaccurate because the decay of  $G_2(t)$  is very small within the experimentally accessible time range. Nonetheless it is clear that the relative decrease of  $\lambda$  with addition of Te is greater than in the low-T region. Although part of this decrease may be ascribed to change of  $\langle v_Q^2 \rangle$ ,  $\tau_c$  must also be decreased by addition of Te. This is in accord with naive expectations, since adding Te is known to weaken the bonding. The "intrinsic"  $\tau_c$  temperature dependence cannot be determined with confidence from the present data. If  $\tau_c$  is activated in this region, the activation energy is of order 0.9 eV for the higher Te-content alloys, approximately the same as the slope of the line which connects the low and high-T "intrinsic" pure Se correlation times in Fig. 2



More accurate data in this region are required before one can draw any reliable conclusions about the temperature dependence of  $\tau_c$  in this region however.

#### ACKNOWLEDGMENTS

We appreciate several fruitful discussions with Professor Melvin Cutler.

This work was supported in part by the Office of Naval Research.

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